Power laws in the dynamic hysteresis of quantum nonlinear photonic resonator

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- optical bistability: semiclassical approach
- optical bistability at the quantum level
- Dynamical hysteresis
- Conclusions & perspectives

Introduction: Kerr model

 $\hbar = 1$

System Hamiltonian:
$$\hat{H}_S = \hbar \omega_c \hat{a}^{\dagger} \hat{a} + \frac{U}{2} \hat{a}^{\dagger} \hat{a}^{\dagger} \hat{a} \hat{a}$$

Coherent drive term: $\hat{H}_P = F(t) e^{-i\omega_p t} \hat{a}^{\dagger} + F^*(t) e^{i\omega_p t} \hat{a}$
Total Hamiltonian: $\hat{H} = \hat{H}_S + \hat{H}_P$

Losses described by Lindblad master equation for density operator:

$$\partial_t \hat{
ho} ~~=~~-i\left[\hat{H},\hat{
ho}
ight]+rac{\gamma}{2}\left(2\hat{a}^\dagger\hat{
ho}\hat{a}-\hat{a}^\dagger\hat{a}\hat{
ho}-\hat{
ho}\hat{a}^\dagger\hat{a}
ight)$$



Review article: I. Carusotto and C. Ciuti – Rev. Mod. Phys. 85, 299 (2013)

Experimental implementations



Optical bistability





H. M. Gibbs, S. L. McCall, and T. N. C. Venkatesan Phys. Rev. Lett. **36**, 1135 (1976)

Optical bistability: semiclassical approach



H. M. Gibbs, S. L. McCall, and T. N. C. Venkatesan Phys. Rev. Lett. **36**, 1135 (1976)

optical bistability at the quantum level

The quantum solution is unique:



P. D. Drummond and D. F. Walls – J. Phys. A: Math. Gen. 13, 725 (1980).

optical bistability at the quantum level



P. D. Drummond and D. F. Walls – J. Phys. A: Math. Gen. 13, 725 (1980).

Bimodal Wigner distribution:



optical bistability at the quantum level



J. Kerckhoff, M. A. Armen, and H. Mabuchi, Opt. Express 19, 24468 (2011).

So far: analysis of the steady-state properties Q: what about time dependence of the experimental sweep?

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Dynamical hysteresis:



Power law behavior of the hysteresis area A:



$$\partial_t \hat{
ho} ~=~ -i \left[\hat{H}, \hat{
ho}
ight] + rac{\gamma}{2} \left(2 \hat{a}^\dagger \hat{
ho} \hat{a} - \hat{a}^\dagger \hat{a} \hat{
ho} - \hat{
ho} \hat{a}^\dagger \hat{a}
ight)$$



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Characteristic timescales
reaction time:
$$\tau_R = -1/Re[\lambda]$$

distance from transition: $\epsilon(t) = F_c - F(t)$
transition rate: $\left|\frac{\dot{\epsilon}(t)}{\epsilon(t)}\right| = \frac{\Delta F}{t_s} \frac{1}{|F_c - F(t)|} = \frac{1}{\tau_s}$ Sweep timescale



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Conclusions & perspectives

- We revealed the dynamical character of the optical bistability hysteresis cycle which is qualitatively different from the semiclassical static hysteresis.
- We showed that the behavior of the Liouvillian gap allows to qualitatively understand the double power law behavior of the hysteresis area.

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What about coupled arrays of photonic resonators?



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